

THE IR STRUCTURE OF CELESTIAL GLUON AMPLITUDES

HERNAN GONZALEZ

IN COLLABORATION WITH FRANCISCO ROJAS

FACULTAD DE ARTES LIBERALES
UNIVERSIDAD ADOLFO IBAÑEZ

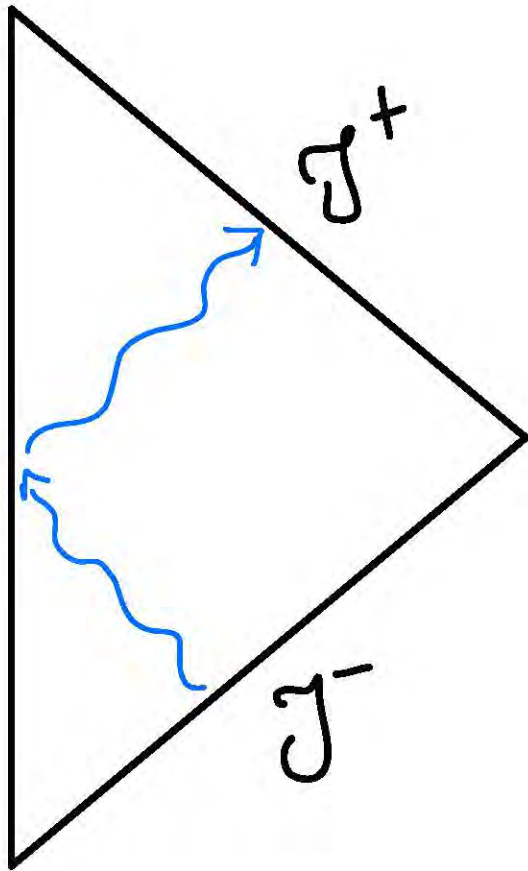
MOTIVATION

Flat Space Holography



SYMMETRIES

MOTIVATION



Flat Space Holography



SYMMETRIES

- GRAVITY + BOUNDARY CONDITIONS AT \mathcal{I}^{\pm}

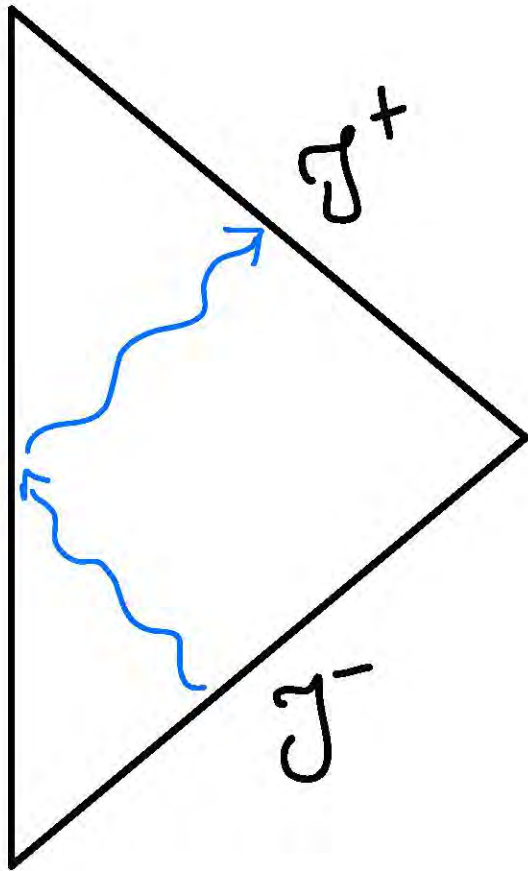
∞ -DIMENSIONAL

SYMMETRIES

(B.M.S.) [1962]

\supset POINCARÉ

MOTIVATION



Flat Space Holography



SYMMETRIES

- GRAVITY + BOUNDARY CONDITIONS AT \mathcal{I}^{\pm}

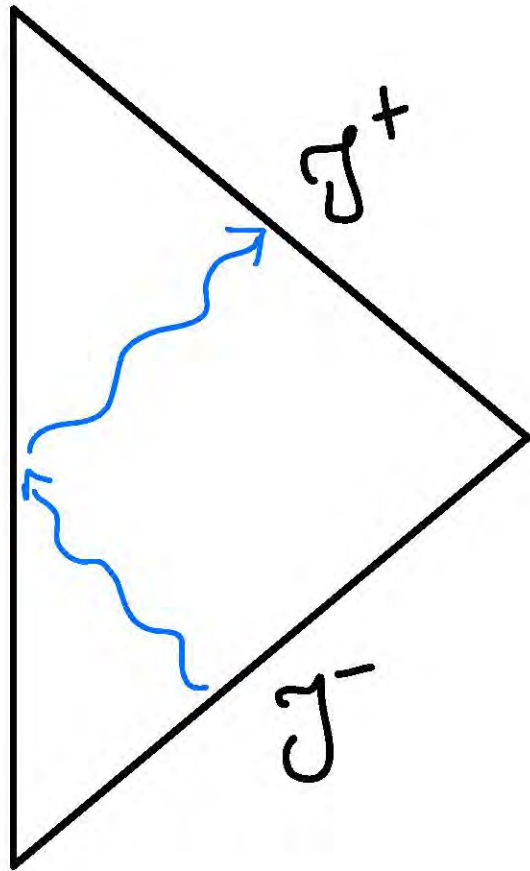
∞ -DIMENSIONAL

SYMMETRIES

(B.M.S.) [1962]

\supset POINCARÉ

MOTIVATION



Flat Space Holography



SYMMETRIES

- GRAVITY + BOUNDARY CONDITIONS AT \mathcal{J}^{\pm}

∞ -DIMENSIONAL

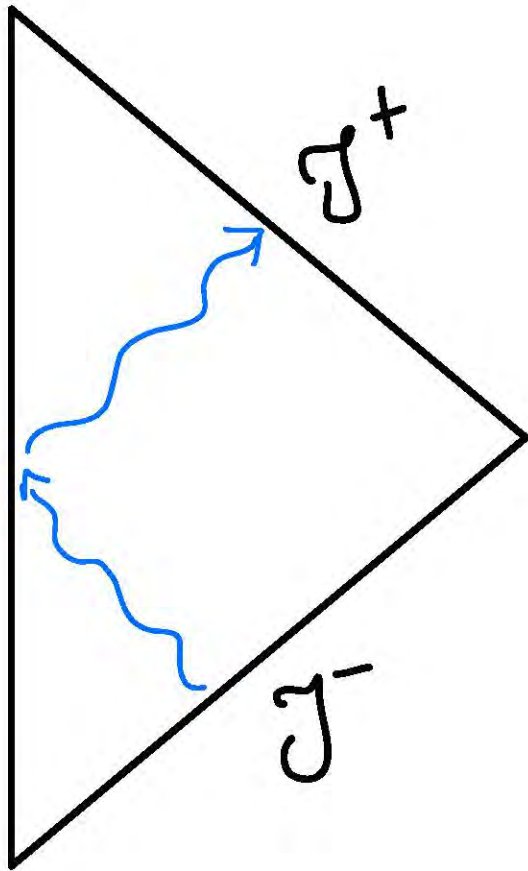
SYMMETRIES

(B.M.S.) [1962]

\supset POINCARÉ

[BARNICH-TROESSART (2010) CAMPIGLIA-LADDHA (2015)]

MOTIVATION



Flat Space Holography



SYMMETRIES

- GRAVITY + BOUNDARY CONDITIONS AT J^{\pm}

∞ -DIMENSIONAL

SYMMETRIES

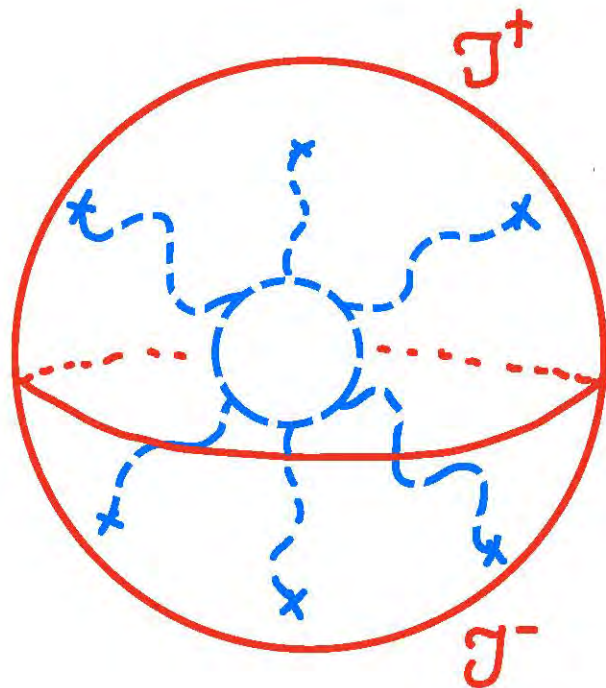
\supset POINCARÉ

(B.M.S.) [1962]

[BARNICH-TROESSART (2010) CAMPIGLIA-LADDHA (2015)]

WHAT ARE THE ROLE OF THESE
SYMMETRIES AT QUANTUM LEVEL?

S-MATRIX FOR MASSLESS PARTICLES



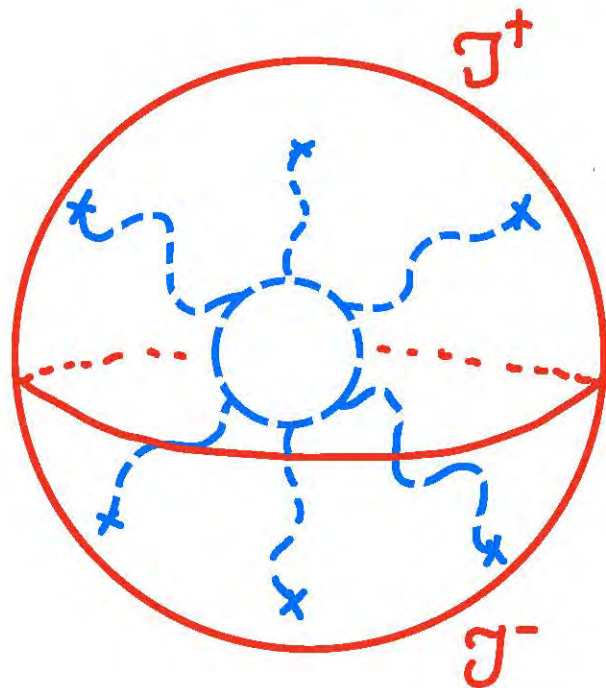
$$\langle \text{OUT} | S | \text{IN} \rangle$$

} }

J^+ J^-

4D BULK / 2D CELESTIAL SPHERE

S-MATRIX FOR MASSLESS PARTICLES



$$\langle \text{OUT} | S | \text{IN} \rangle$$

} }

J^+ J^-

4D BULK / 2D CELESTIAL SPHERE

WHAT ARE THE PROPERTIES OF THIS 2D THEORY?

OUTLINE

- Mellin Basis & Symmetries
- Celestial Factorization in IR Regulated Amplitudes
- CFT for IR divergences
- Conclusions

A NEW BASIS

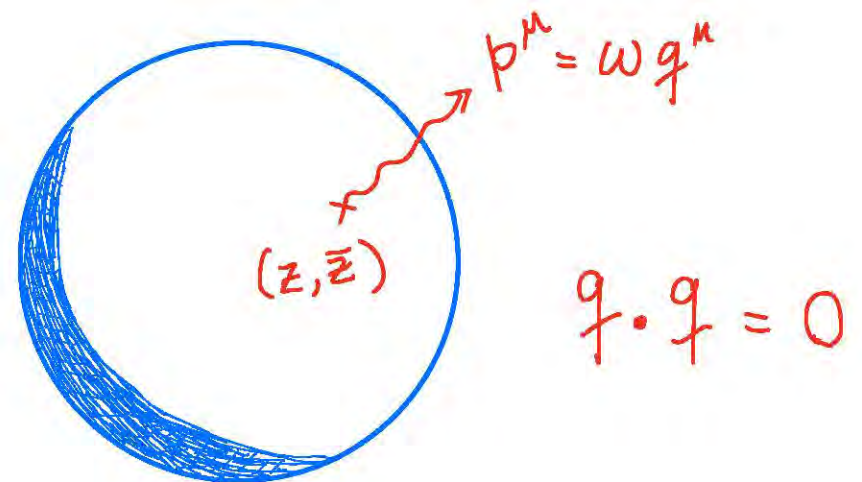
[PASTERSKI & SHAO (2017)]

- USUAL BASIS \rightsquigarrow PLANES WAVES $\Psi_p(x) = e^{i p \cdot x}$

A NEW BASIS

[PASTERSKI & SHAO (2017)]

- USUAL BASIS \rightsquigarrow PLANES WAVES $\Psi_p(x) = e^{i p \cdot x}$
- INSERTIONS ON CS ARE PARAMETERIZED BY MOMENTUM'S PARTICLE:

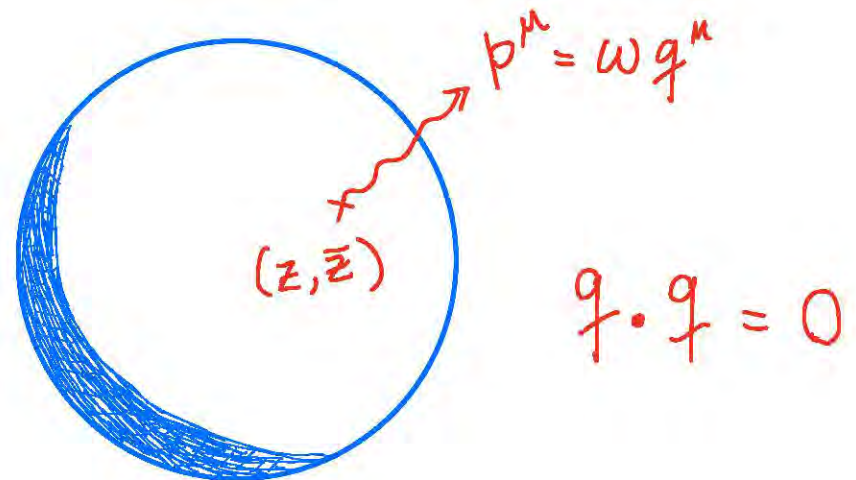


A NEW BASIS

[PASTERSKI & SHAO (2017)]

- USUAL BASIS \rightsquigarrow PLANES WAVES $\Psi_p(x) = e^{i p \cdot x}$
- INSERTIONS ON CS ARE PARAMETERIZED BY MOMENTUM'S PARTICLE:

$$\mathcal{P}^M = \frac{\omega}{2} (1 + |z|^2, z + \bar{z}, -i(z - \bar{z}), 1 - |z|^2) = \omega q^M$$



A NEW BASIS

[PASTERSKI & SHAO (2017)]

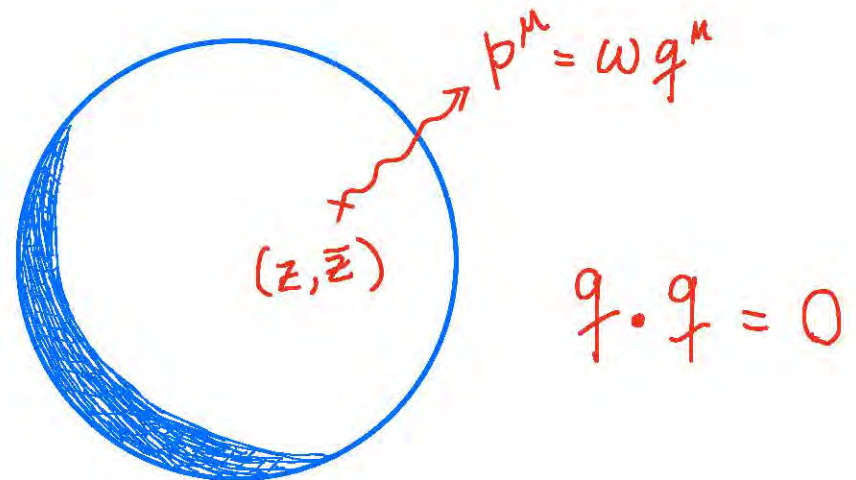
- USUAL BASIS \rightsquigarrow PLANES WAVES $\Psi_p(x) = e^{i p \cdot x}$
- INSERTIONS ON CS ARE PARAMETERIZED BY MOMENTUM'S PARTICLE:

$$P^M = \frac{\omega}{2} (1 + |z|^2, z + \bar{z}, -i(z - \bar{z}), 1 - |z|^2) = \omega q^M$$

$$z \rightarrow z' = \frac{az + b}{cz + d}$$

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{C})$$

GLOBAL
CONFORMAL
SYMMETRY
IN 2D



A NEW BASIS

[PASTERSKI & SHAO (2017)]

○ CHANGE OF BASIS USING MELLIN TRANSFORM

PLANE WAVE : $\exp(i\omega q \cdot x)$
(SCALAR)

ENERGY
↑
POSITION
ON CS

NEW STATE : $O_{\Delta}(z; x)$

A NEW BASIS

[PASTERSKI & SHAO (2017)]

○ CHANGE OF BASIS USING MELLIN TRANSFORM

PLANE WAVE :
(SCALAR)

$$\exp(i\omega q \cdot x)$$

↑ ENERGY
↓ POSITION ON CS

NEW STATE :

$$O_{\Delta}(z; x) = \int_0^{\infty} \frac{d\omega}{\omega} \omega^{\Delta} e^{i\omega q \cdot x}$$

A NEW BASIS

[PASTERSKI & SHAO (2017)]

○ CHANGE OF BASIS USING MELLIN TRANSFORM

PLANE WAVE :
(SCALAR)

$$\exp(i\omega q \cdot x)$$

↑ ENERGY
↓ POSITION ON CS

NEW STATE :

$$O_{\Delta}(z; x) = \int_0^{\infty} \frac{d\omega}{\omega} \omega^{\Delta} e^{i\omega q \cdot x}$$

○ SUMMING OVER ALL ENERGY SCALES

A NEW BASIS

[PASTERSKI & SHAO (2017)]

○ CHANGE OF BASIS USING MELLIN TRANSFORM

PLANE WAVE :
(SCALAR)

$$\exp(i\omega q \cdot x)$$

↑ ENERGY
↓ POSITION ON CS

NEW STATE :

$$O_{\Delta}(z; x) = \int_0^{\infty} \frac{d\omega}{\omega} \omega^{\Delta} e^{i\omega q \cdot x}$$

- SUMMING OVER ALL ENERGY SCALES
- EXCHANGING ENERGY (ω) BY SCALING DIMENSION (Δ)

CONFORMAL BASIS

- UNDER $SL(2, \mathbb{C})$: TRANSFORMS AS A PRIMARY FIELD

$$\mathcal{O}_{\Delta}(\Lambda^{\mu}_{\nu} x^{\nu}, \frac{az+b}{cz+d}) = |cz+d|^{-\Delta} \mathcal{O}_{\Delta}(x^{\mu}; z)$$

CONFORMAL BASIS

○ UNDER $SL(2, \mathbb{C})$: TRANSFORMS AS A PRIMARY FIELD

$$\mathcal{O}_{\Delta} \left(\Lambda^{\mu}_{\nu} x^{\nu}, \frac{az+b}{cz+d} \right) = |cz+d|^{\Delta} \mathcal{O}_{\Delta} (x^{\mu}; z)$$

$$\downarrow$$
$$\begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

↳ WITH SPIN = 0.
(SCALAR)

CONFORMAL BASIS

- UNDER $SL(2, \mathbb{C})$: TRANSFORMS AS A PRIMARY FIELD

$$\mathcal{O}_{\Delta} \left(\Lambda^{\mu}_{\nu} x^{\nu}, \frac{az+b}{cz+d} \right) = |cz+d|^{\Delta} \mathcal{O}_{\Delta} (x^{\mu}; z)$$

\downarrow
 $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$

- DELTA - NORMALIZABLE $(\Psi_p(x), \Psi_{p'}(x)) = (2\pi)^3 (2p^0) \delta^{(3)}(p-p')$

SAME NORMALIZATION ON CONFORMAL BASIS $\rightarrow \Delta = 1 + i\lambda$

CONFORMAL BASIS

- UNDER $SL(2, \mathbb{C})$: TRANSFORMS AS A PRIMARY FIELD

$$\mathcal{O}_{\Delta} \left(\Lambda^{\mu}_{\nu} x^{\nu}, \frac{az+b}{cz+d} \right) = |cz+d|^{\Delta} \mathcal{O}_{\Delta} (x^{\mu}; z)$$

\downarrow
 $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$

- DELTA - NORMALIZABLE $(\Psi_p(x), \Psi_{p'}(x)) = (2\pi)^3 (2p^0) \delta^{(3)}(p-p')$

CONFORMAL BASIS

- UNDER $SL(2, \mathbb{C})$: TRANSFORMS AS A PRIMARY FIELD

$$\mathcal{O}_{\Delta} \left(\Lambda^{\mu}_{\nu} x^{\nu}, \frac{az+b}{cz+d} \right) = |cz+d|^{\Delta} \mathcal{O}_{\Delta} (x^{\mu}; z)$$

\downarrow
 $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$

- DELTA - NORMALIZABLE $(\Psi_p(x), \Psi_{p'}(x)) = (2\pi)^3 (2p^0) \delta^{(3)}(p-p')$

SAME NORMALIZATION ON CONFORMAL BASIS $\rightarrow \Delta = 1 + i\epsilon$



DIFFERENT KIND OF CFT

CELESTIAL AMPLITUDES

[PASTERSKI, SHAO (2017)]

[PASTERSKI, SHAO, STROMINGER (2017)]

○ IN MOMENTUM BASIS:

$$A_n(\{p_i, \epsilon_j\}) = i(2\pi)^4 \delta(p_1 + p_2 - \sum_{k=3}^n p_k) A(\{p_i, \epsilon_i\})$$

CELESTIAL AMPLITUDES

[PASTERSKI, SHAO (2017)]

[PASTERSKI, SHAO, STROMINGER (2017)]

○ IN MOMENTUM BASIS:

$$A_n(\{p_i, \epsilon_j\}) = i(2\pi)^4 \delta(p_1 + p_2 - \sum_{k=3}^n p_k) A(\{p_i, \epsilon_i\})$$

○ USING MELLIN TRANSFORM

$$\tilde{A}_{\{\Delta_k, l_k\}}(\{z_k, \bar{z}_k\}) = \prod_{k=1}^n \int_0^\infty \frac{d\omega_k}{\omega_k} \omega_k^{\Delta_k - 1} A_n(\{p_i, \epsilon_j\})$$

CELESTIAL AMPLITUDES

[PASTERSKI, SHAO (2017)]

[PASTERSKI, SHAO, STROMINGER (2017)]

- IN MOMENTUM BASIS:

$$A_n(\{p_i, \epsilon_j\}) = i(2\pi)^4 \delta(p_1 + p_2 - \sum_{k=3}^n p_k) A(\{p_i, \epsilon_i\})$$

- USING MELLIN TRANSFORM

$$\tilde{A}_{\{\Delta_k, l_k\}}(\{z_k, \bar{z}_k\}) = \underbrace{\prod_{k=1}^n \int_0^\infty \frac{d\omega_k}{\omega_k} \omega_k^{\Delta_k - 1}}_{\text{INTEGRATION OVER EACH EXTERNAL LEG}} A_n(\{p_i, \epsilon_j\})$$

INTEGRATION OVER EACH
EXTERNAL LEG

CELESTIAL AMPLITUDES

[PASTERSKI, SHAO (2017)]

[PASTERSKI, SHAO, STROMINGER (2017)]

○ IN MOMENTUM BASIS:

$$A_n(\{p_i, \epsilon_j\}) = i(2\pi)^4 \delta(p_1 + p_2 - \sum_{k=3}^n p_k) A(\{p_i, \epsilon_i\})$$

○ USING MELLIN TRANSFORM

$$\tilde{A}_{\{\Delta_k, l_k\}}(\{z_k, \bar{z}_k\}) = \underbrace{\prod_{k=1}^n \int_0^\infty \frac{d\omega_k}{\omega_k} \omega_k^{\Delta_k - 1}}_{\text{INTEGRATION OVER EACH EXTERNAL LEG}} A_n(\{p_i, \epsilon_j\})$$

INTEGRATION OVER EACH
EXTERNAL LEG

CELESTIAL AMPLITUDES

[PASTERSKI, SHAO (2017)]

[PASTERSKI, SHAO, STROMINGER (2017)]

- IN MOMENTUM BASIS:

$$A_n(\{p_i, \epsilon_j\}) = i(2\pi)^4 \delta(p_1 + p_2 - \sum_{k=3}^n p_k) A(\{p_i, \epsilon_i\})$$

- USING MELLIN TRANSFORM

$$\tilde{A}_{\{\Delta_k, l_k\}}(\{z_k, \bar{z}_k\}) = \underbrace{\prod_{k=1}^n \int_0^\infty \frac{d\omega_k}{\omega_k} \omega_k^{\Delta_k - 1}}_{\text{INTEGRATION OVER EACH EXTERNAL LEG}} A_n(\{p_i, \epsilon_j\})$$

INTEGRATION OVER EACH
EXTERNAL LEG

- \tilde{A} INVOLVES ARBITRARILY HIGH ENERGY PARTICLES (ω_k)

CELESTIAL = CORRELATION FUNCTION
AMPLITUDE = 2D - CFT

$$\tilde{A}_{\{\Delta_k, l_k\}}(\{z_k, \bar{z}_k\}) = \langle \mathcal{O}_{\Delta_1 l_1}(z_1, \bar{z}_1) \dots \mathcal{O}_{\Delta_n l_n}(z_n, \bar{z}_n) \rangle$$

TRANSFORMATION PROPERTIES

○ UNDER $SL(2, \mathbb{C})$ LORENTZ TRANSFORMATION

$$\tilde{A}_{\{\Delta_k, l_k\}} \left(\left\{ \frac{az_k + b}{cz_k + d}, \frac{\bar{a}\bar{z}_k + \bar{b}}{\bar{c}\bar{z}_k + \bar{d}} \right\} \right) = \prod_{k=1}^n \left((cz_k + d)^{\Delta_k + l_k} (\bar{c}\bar{z}_k + \bar{d})^{\Delta_k - l_k} \right) \tilde{A}_{\{\Delta_k, l_k\}} (\{z_k, \bar{z}_k\})$$

$$h_k = \frac{\Delta_k + l_k}{2}$$

$$\bar{h}_k = \frac{\Delta_k - l_k}{2}$$

FOR EACH EXTERNAL PARTICLE

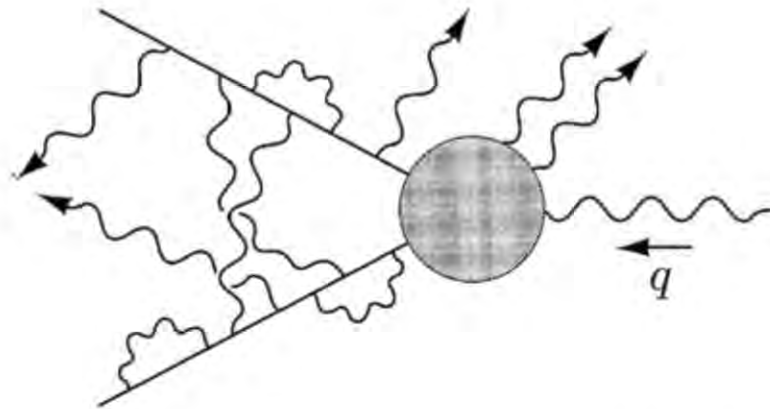
l_k : HELICITY

GLUON AMPLITUDES AT LOOP LEVEL

- MELLIN AMPLITUDES COMBINES HIGH AND LOW ENERGY STATES.
- THEORIES WITH MASSLESS EXCITATIONS (PHOTONS, GLUONS, ...)
SUFFERS FROM IR RADIATION AT LOOP LEVEL.

GLUON AMPLITUDES AT LOOP LEVEL

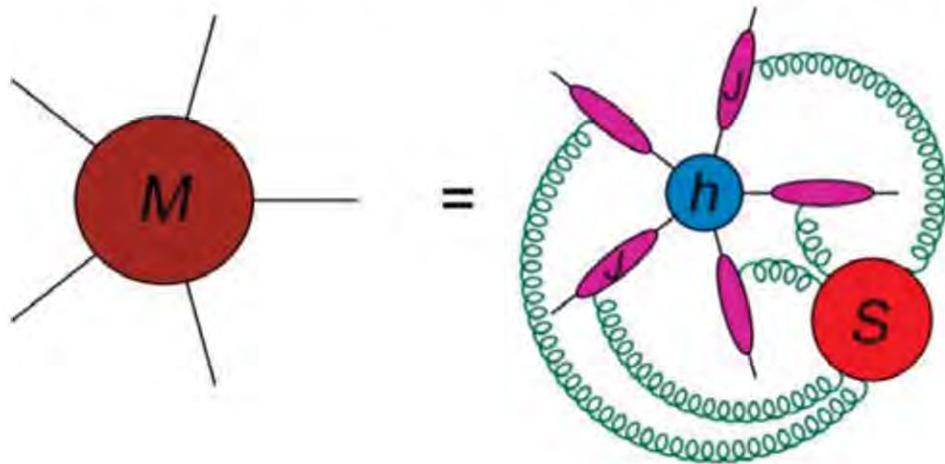
- MELLIN AMPLITUDES COMBINES HIGH AND LOW ENERGY STATES.
- THEORIES WITH MASSLESS EXCITATIONS (PHOTONS, GLUONS, ...)
SUFFERS FROM IR RADIATION AT LOOP LEVEL.



IR PROBLEM : ∞ INTERNAL EXCHANGES
WITH VERY LOW ENERGY

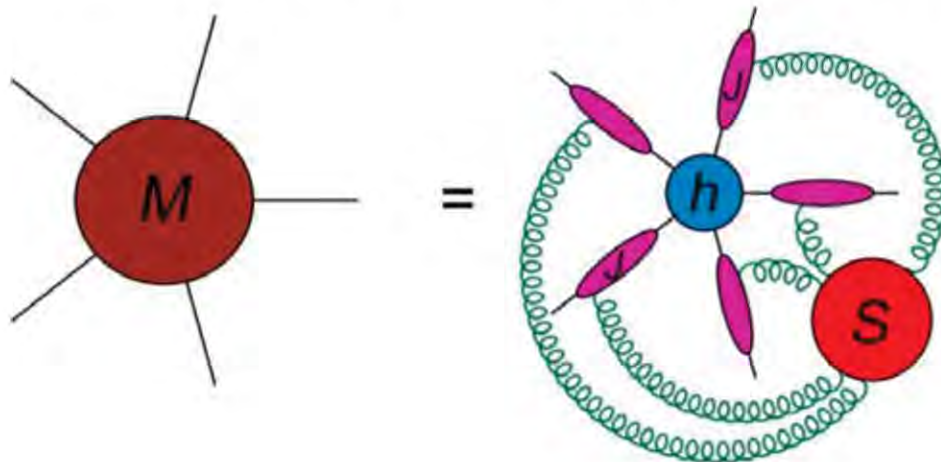
GLUON AMPLITUDES AT LOOP LEVEL

- MELLIN AMPLITUDES COMBINES HIGH AND LOW ENERGY STATES.
- THEORIES WITH MASSLESS EXCITATIONS (PHOTONS, GLUONS, ...) SUFFERS FROM IR RADIATION AT LOOP LEVEL.
- FORTUNATELY, IT IS KNOWN HOW TO DEAL WITH THEM CONSIDERING ALL LOOP RESSUMATIONS (SOFT LIMIT)



GLUON AMPLITUDES AT LOOP LEVEL

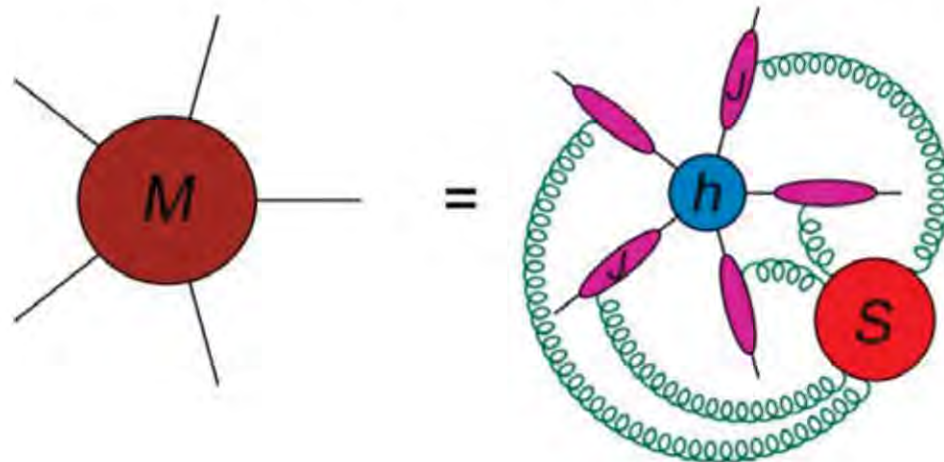
- MELLIN AMPLITUDES COMBINES HIGH AND LOW ENERGY STATES.
- THEORIES WITH MASSLESS EXCITATIONS (PHOTONS, GLUONS, ...) SUFFERS FROM IR RADIATION AT LOOP LEVEL.
- FORTUNATELY, IT IS KNOWN HOW TO DEAL WITH THEM CONSIDERING ALL LOOP RESSUMATIONS (SOFT LIMIT)



$$M = \underbrace{\left(\prod_{i=1}^n J_i \right) S}_{\text{IR RESSUMATION}} \mathcal{H}$$

GLUON AMPLITUDES AT LOOP LEVEL

- MELLIN AMPLITUDES COMBINES HIGH AND LOW ENERGY STATES.
- THEORIES WITH MASSLESS EXCITATIONS (PHOTONS, GLUONS, ...) SUFFERS FROM IR RADIATION AT LOOP LEVEL.
- FORTUNATELY, IT IS KNOWN HOW TO DEAL WITH THEM CONSIDERING ALL LOOP RESSUMATIONS (SOFT LIMIT)



$$M = \underbrace{\left(\prod_{i=1}^n J_i \right)}_{\text{IR RESSUMATION}} S \overset{\text{HARD PIECE}}{\mathcal{H}}$$

CELESTIAL n -POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY

CELESTIAL n -POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY
- FACTORIZATION BETWEEN SOFT AND HARD DEGREES OF FREEDOM

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY
- FACTORIZATION BETWEEN SOFT AND HARD DEGREES OF FREEDOM

$$A_n(\{p_i, \epsilon_i\}) = Z(\mu^2, p_i \cdot p_j) A_{\text{finite}}(\{p_i, \epsilon_i\})$$

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY
- FACTORIZATION BETWEEN SOFT AND HARD DEGREES OF FREEDOM

$$A_n(\{p_i, \epsilon_i\}) = \mathcal{Z}(\mu^2, p_i \cdot p_j) A_{\text{finite}}(\{p_i, \epsilon_i\})$$

↗ IR-DIVERGENT OPERATOR " $(\prod_{i=1}^n J_i) S$ "

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY
- FACTORIZATION BETWEEN SOFT AND HARD DEGREES OF FREEDOM

$$A_n(\{p_i, \epsilon_i\}) = \int (\mu^2, p_i \cdot p_j) \underbrace{A_{\text{finite}}(\{p_i, \epsilon_i\})}_{\text{VECTOR DECOMPOSED IN COLOR SPACE}}$$

↗ IR-DIVERGENT OPERATOR $\left(\prod_{i=1}^n J_i\right) S$

- DIMENSIONAL REGULARIZATION $D=4-2\epsilon$ $\epsilon < 0$ FOR INTERNAL INTERACTIONS

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- WE CONSIDER A PROCESS INVOLVING n EXTERNAL GLUONS WITH HELICITIES $\lambda_i = \pm$ AND $SU(N)$ INTERNAL SYMMETRY
- FACTORIZATION BETWEEN SOFT AND HARD DEGREES OF FREEDOM

$$A_n(\{p_i, \epsilon_i\}) = Z(\mu^2, p_i \cdot p_j) \underbrace{A_{\text{finite}}(\{p_i, \epsilon_i\})}_{\text{VECTOR DECOMPOSED IN COLOR SPACE}}$$

↗ IR-DIVERGENT OPERATOR " $(\prod_{i=1}^n J_i) S$ "

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

n-POINT AMPLITUDES

$$M_n(p_1, \dots, p_n) = \underbrace{Z_n}_{\substack{\text{DIVERGENT} \\ \text{FOR } \Lambda \rightarrow 0}} \times \underbrace{A_n}_{\text{FINITE}}$$

EVOLUTION EQUATIONS

$$\frac{d}{d \log \mu^2} Z_n = -\Gamma_n Z_n$$

↓ ENERGY SCALE

↘ ANOMALOUS DIMENSION

$$Z(\mu^2, p_i \cdot p_j) = \mathcal{P} \exp \left\{ -\frac{1}{2} \int_0^{\mu^2} \frac{d\lambda^2}{\lambda^2} \Gamma_n(p_i \cdot p_j, \lambda^2) \right\}$$

CELESTIAL n -POINT GLUON AMPLITUDE (ALL LOOPS)

MAIN OBJECTIVE :

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

MAIN OBJECTIVE:

CELESTIAL CFT BEHIND CORRELATOR

$$Z_n(\mu^2, p_i \cdot p_j)$$

CELESTIAL n -POINT GLUON AMPLITUDE (ALL LOOPS)

MAIN OBJECTIVE:

CELESTIAL CFT BEHIND CORRELATOR

$$Z_n(\mu^2, p_i \cdot p_j)$$

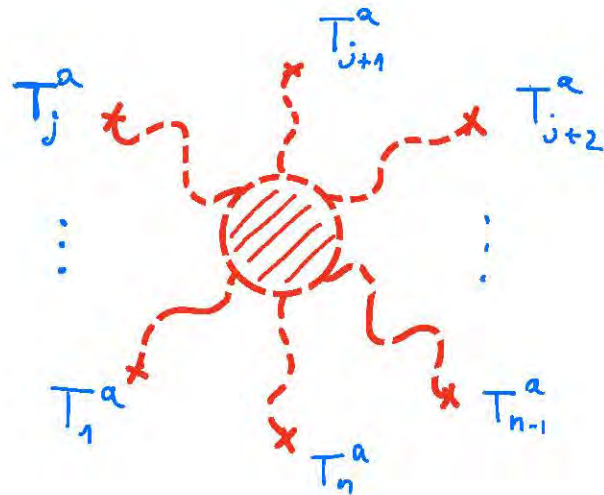
1) $N \gg 1$

2) FINITE N
TRUNCATED IN α_s

[MAGNEA 2021]

COLORED
FREE SCALAR

○ SU(N) GLUON AMPLITUDES (COLOR SPACE FORMALISM CATANI [1998])



T_i^a : COLOR CHARGE OPERATOR
IN ADJOINT REPRESENTATION

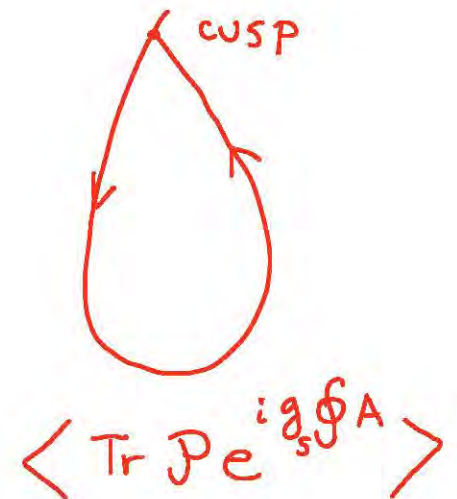
$$\sum_{i=1}^n T_i^a = 0 : \text{COLOR CHARGE CONSERVATION}$$

GARDI-MAGNEA
BECHER-NEUBERT
[2009]

○ IMPORTANT CONSTRAINT IN QCD AMPLITUDES:

$$\sum_{j \neq i} \frac{\partial}{\partial \log p_{ij}} \Gamma_n^{(\text{SOFT})} = \gamma_{\text{cusp}}^i(\alpha_s) \mathbb{1}$$

p_{ij} invariant under $p_i \rightarrow \lambda_i p_i$



CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- MELLIN AMPLITUDE FACTORIZES AS WELL:

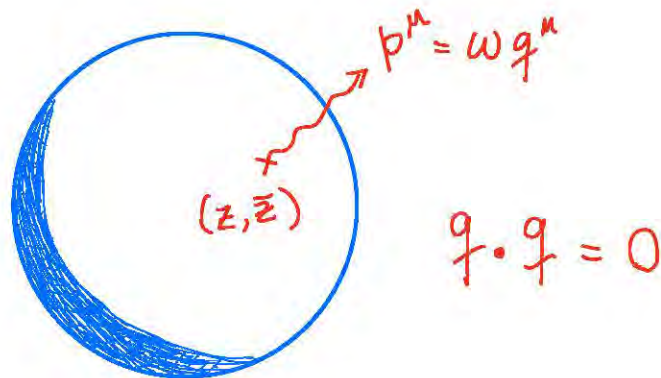
$$\tilde{\mathcal{A}}_n^{\text{ALL-LOOPS}} = \mathcal{Z}_n(\mu^2, |z_{ij}|^2) \tilde{\mathcal{A}}_{\text{finite}} \left(h_i^? = h_i - \frac{1}{4}K \right)$$

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- MELLIN AMPLITUDE FACTORIZES AS WELL:

$$\tilde{\mathcal{A}}_n^{\text{ALL-LOOPS}} = \mathcal{Z}_n(\mu^2, |z_{ij}|^2) \tilde{\mathcal{A}}_{\text{finite}} \left(h_i^? = h_i - \frac{1}{4}K \right)$$

RECALL:

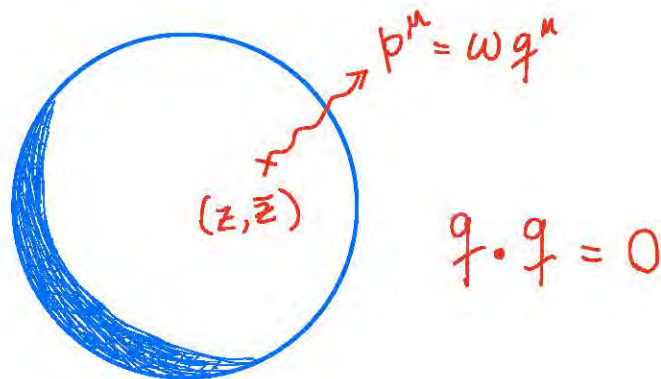


CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- MELLIN AMPLITUDE FACTORIZES AS WELL:

$$\tilde{\mathcal{A}}_n^{\text{ALL-LOOPS}} = \mathcal{Z}_n(\mu^2, |z_{ij}|^2) \tilde{\mathcal{A}}_{\text{finite}} \left(h_i^? = h_i - \frac{1}{4}K \right)$$

RECALL:



GARDI-MAGNEA
BECHER-NEUBERT
[2009]

$$\sum_{j \neq i} \frac{\partial}{\partial \log \rho_{ij}} \Gamma_n^{(\text{SOFT})} = \gamma_{\text{CUSP}}^i(\alpha_s) \mathbb{1}$$

CELESTIAL n -POINT GLUON AMPLITUDE (ALL LOOPS)

- MELLIN AMPLITUDE FACTORIZES AS WELL:

$$\tilde{\mathcal{A}}_n^{\text{ALL-LOOPS}} = \mathcal{Z}_n(\mu^2, |z_{ij}|^2) \tilde{\mathcal{A}}_{\text{finite}} \quad (h_i = h_i - \frac{1}{4}K)$$

- INTERESTING PHYSICS ENCODED IN \mathcal{Z}_n SPECIFIC THEORY?

IT CAN BE WRITTEN IN TERMS OF n -POINT CORRELATORS OF VERTEX OPERATORS $\mathcal{V}_i(z, \bar{z})$ WITH CONFORMAL DIMENSION $\frac{1}{4}K$

CELESTIAL n-POINT GLUON AMPLITUDE (ALL LOOPS)

- MELLIN AMPLITUDE FACTORIZES AS WELL:

$$\tilde{\mathcal{A}}_n^{\text{ALL-LOOPS}} = \mathcal{Z}_n(\mu^2, |z_{ij}|^2) \tilde{\mathcal{A}}_{\text{finite}} \quad (h_i = h_i - \frac{1}{4}K)$$

- INTERESTING PHYSICS ENCODED IN \mathcal{Z}_n SPECIFIC THEORY?

IT CAN BE WRITTEN IN TERMS OF n-POINT CORRELATORS OF VERTEX OPERATORS $\mathcal{V}_i(z, \bar{z})$ WITH CONFORMAL DIMENSION $\frac{1}{4}K$

$$\mathcal{Z}_n = \langle \mathcal{V}_{T_1} \otimes \mathcal{V}_{T_2} \otimes \dots \otimes \mathcal{V}_{T_n} \rangle$$

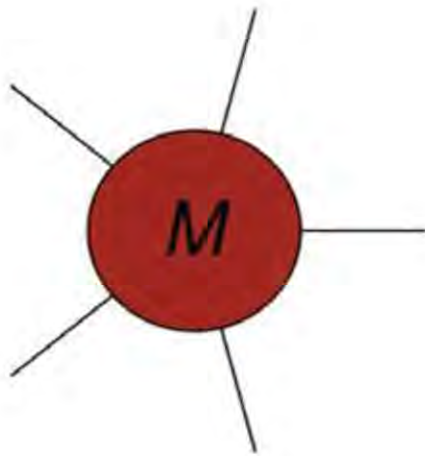
2D - CFT CORRELATOR
GARDI, SMILLIE [2011, 2013]
WHITE

$$K = \frac{1}{2} \int_0^{\mu^2} \frac{d\lambda^2}{\lambda^2} \gamma_{\text{CUSP}}(\alpha_s(\lambda^2))$$

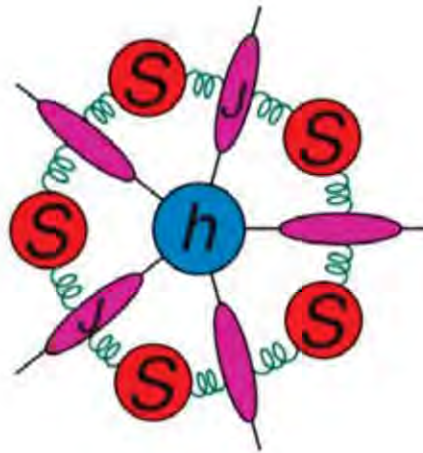
DIVERGES WITH
THE REGULATOR ϵ

LARGE N

$$Z_n = \exp \left(-\frac{Nk}{4} \sum_{i < j} \overbrace{(\delta_{i,j+1} + \delta_{j,i+1})}^{\text{NEAREST NEIGHBOUR INTERACTIONS}} \log(|z_{ij}|^2) \right)$$



=



- ONLY PLANAR DIAGRAMS CONTRIBUTE
- INTERACTIONS BETWEEN NEIGHBOURING LEGS

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\mathcal{K} = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\mathcal{K} = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

COULOMB GAS DESCRIPTION :

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\kappa = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

COULOMB GAS DESCRIPTION: $\langle \phi^a(z) \phi^b(0) \rangle = -\kappa \delta^{ab} \log(|z_{ij}|^2)$

↑
FREE BOSON WITH A COLOR INDEX

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\kappa = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

COULOMB GAS DESCRIPTION: $\langle \phi^a(z) \phi^b(0) \rangle = -\kappa \delta^{ab} \log(|z_{ij}|^2)$

↑
FREE BOSON WITH A COLOR INDEX

$$Z_n = \langle :e^{i\tau_1 \cdot \phi}: :e^{i\tau_2 \cdot \phi}: \dots :e^{i\tau_n \cdot \phi}: \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\kappa = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

COULOMB GAS DESCRIPTION: $\langle \phi^a(z) \phi^b(0) \rangle = -\kappa \delta^{ab} \log(|z_{ij}|^2)$

↑
FREE BOSON WITH A COLOR INDEX

$$Z_n = \langle :e^{i\tau_1 \cdot \phi} : :e^{i\tau_2 \cdot \phi} : \dots :e^{i\tau_n \cdot \phi} : \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

NO LONGER
AN OPERATOR $\rightarrow T_i^a$

LARGE N

$$Z_n = \exp \left(-\kappa \sum_{i < j} (\delta_{i,j+1} + \delta_{j,i+1}) \log(|z_{ij}|^2) \right) \quad \begin{array}{l} 4\kappa = NK = \text{fixed} \\ N \rightarrow \infty \end{array}$$

COULOMB GAS DESCRIPTION: $\langle \phi^a(z) \phi^b(0) \rangle = -\kappa \delta^{ab} \log(|z_{ij}|^2)$
 \uparrow
 FREE BOSON WITH A COLOR INDEX

$$Z_n = \langle :e^{i\tau_1 \cdot \phi} : :e^{i\tau_2 \cdot \phi} : \dots :e^{i\tau_n \cdot \phi} : \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

NO LONGER AN OPERATOR \rightarrow

$$T_i^a = \underbrace{(0, \dots, \overset{\text{POSITION } i}{\downarrow} 1, \overset{\text{POSITION } i+1}{\downarrow} -1, 0, \dots, 0)}_{N^2 \rightarrow \infty \text{ COMPTES}} \rightsquigarrow \sum_{i=1}^n T_i^a = 0$$

LARGE N

$$Z_n = \langle :e^{i\tau_1 \cdot \phi} : :e^{i\tau_2 \cdot \phi} : \dots :e^{i\tau_n \cdot \phi} : \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

LARGE N

$$Z_n = \langle :e^{i\tau_1 \cdot \phi} : :e^{i\tau_2 \cdot \phi} : \dots :e^{i\tau_n \cdot \phi} : \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

$$V_{\tau_i} = e^{i\tau_i \cdot \phi}$$

VERTEX OPERATOR



$$I(\phi^a) = \frac{1}{8\pi\kappa} \int d^2z \bar{\partial} \phi^a \partial \phi^a$$

EFFECTIVE ACTION

LARGE N

$$Z_n = \langle :e^{i\tau_1 \cdot \phi} : :e^{i\tau_2 \cdot \phi} : \dots :e^{i\tau_n \cdot \phi} : \rangle = \exp \left[- \sum_{i < j} \langle \phi \cdot \tau_i, \phi \cdot \tau_j \rangle \right]$$

$$V_{\tau_i} = e^{i\tau_i \cdot \phi}$$

VERTEX OPERATOR

FINITE N

Γ_n AT FINITE N AND TWO-LOOP TRUNCATION

$$Z_n = \exp\left(\frac{K}{2} \sum_{i < j} T_i \cdot T_j \log(|z_{ij}|^2)\right)$$

ALSO ADMITS A COLORED FREE SCALAR REP.

[MAGNEA 2021]

FINITE N : BEYOND DIPOLAR APPROXIMATION

$$Z_n = \mathcal{P} \exp \left\{ -\frac{1}{2} \int_0^{\mu^2} \frac{d\lambda^2}{\lambda^2} \Gamma_n (|z_{ij}|^2, \lambda^2) \right\}$$

Γ_n ANOMALOUS DIMENSION CONTRIBUTIONS

Kinematic Dependence	Tensor Structure	Loop order	Reference
$\frac{z_{ij} z_{kl}}{z_{il} z_{jk}}$ <p>SL(2,C) INVARIANT RATIOS</p>	$f_{abc} f^c_{\quad d} \{T_i^a, T_j^b\} \{T_j^c, T_k^d\}$	3	ALMELID, DUHR, GARDI, MCLEOD WHITE [2016] [2017]
$\log(z_{ij} ^2)$	$d_{abcd} \text{tr}(T_i^a T_j^b T_k^c T_l^d)$	4	BECHER NEUBERT [2019]

FINITE N : BEYOND DIPOLAR APPROXIMATION

- CCFT CONTAINING NON-ABELIAN STRUCTURE

FINITE N : BEYOND DIPOLAR APPROXIMATION

- CCFT CONTAINING NON-ABELIAN STRUCTURE

↳ BOUNDARY ACTIONS FROM
YM SYMPLECTIC SYMMETRIES?

FINITE N : BEYOND DIPOLAR APPROXIMATION

- CCFT CONTAINING NON-ABELIAN STRUCTURE

↳ BOUNDARY ACTIONS FROM
YM SYMPLECTIC SYMMETRIES?

- 3 AND 4 LOOPS STRUCTURES?

FINITE N : BEYOND DIPOLAR APPROXIMATION

- CCFT CONTAINING NON-ABELIAN STRUCTURE

↳ BOUNDARY ACTIONS FROM
YM SYMPLECTIC SYMMETRIES?

- 3 AND 4 LOOPS STRUCTURES?

- IR SAFE CELESTIAL AMPLITUDES $\mathcal{O}_{(\Delta, l)}^{\text{BARE}} = \sum_i \mathcal{O}_{(\Delta, l)}^{\text{DRESSED}}$

FINITE N : BEYOND DIPOLAR APPROXIMATION

- CCFT CONTAINING NON-ABELIAN STRUCTURE

↳ BOUNDARY ACTIONS FROM
YM SYMPLECTIC SYMMETRIES?

- 3 AND 4 LOOPS STRUCTURES?

- IR SAFE CELESTIAL AMPLITUDES $\mathcal{O}_{(\Delta, l)}^{\text{BARE}} = \sum_i \mathcal{O}_{(\Delta, l)}^{\text{DRESSED}}$

↳ FADDEEV - KULISH CONSTRUCTION

CONCLUSION

- NEW KIND OF HOLOGRAPHY:

CELESTIAL AMPLITUDES \longleftrightarrow 2D-CFT

CONCLUSION

- NEW KIND OF HOLOGRAPHY:

CELESTIAL AMPLITUDES \longleftrightarrow 2D-CFT

- $SU(N)$ GAUGE THEORY, CELESTIAL AMPLITUDE SPLITS INTO HARD AND INFRARED DIVERGENT PIECES.

CONCLUSION

- NEW KIND OF HOLOGRAPHY:

CELESTIAL AMPLITUDES \longleftrightarrow 2D-CFT

- SU(N) GAUGE THEORY, CELESTIAL AMPLITUDE SPLITS INTO HARD AND INFRARED DIVERGENT PIECES,
- CFT IDENTIFICATION FOR INFRARED DIVERGENCES

$$\mathbb{Z}_n = \langle V_{T_1} \otimes V_{T_2} \otimes \dots \otimes V_{T_n} \rangle$$

- LARGE N
- FINITE N
- TRUNCATED

$$V_i = e^{i T_i \cdot \phi} \text{ PRIMARIES } \phi^a \rightsquigarrow \text{ COLORED SCALAR}$$

THANKS!

FINITE N

WZW SATISFIES KNIZHNIK-ZAMOLODCHIKOV EQUATIONS

$$\left(\frac{\partial}{\partial z_i} - \frac{k}{2} \sum_{i \neq j} \frac{T_i T_j}{z_i - z_j} \right) Z_n = 0$$

IN OUR CASE THERE IS AN EXTRA TERM ON THE RIGHT HAND SIDE

Z_n NOT COMPATIBLE WITH WZW

CONFORMAL BASIS

[STIEBERGER & TAYLOR
2019]

○ UNDER SPACE-TIME TRANSLATION : $\hat{P}^\mu = \hat{\omega} q^\mu(z, \bar{z})$

$$\delta_{\hat{p}^\mu} O_\Delta(z; x) = q^\mu \int_0^\infty \frac{d\omega}{\omega} \omega^{(\Delta+1)} e^{i\omega q \cdot x}$$
$$\propto O_{\Delta+1}(z; x)$$

IN PARTICULAR $\hat{P}^+ \equiv \hat{P}^0 + \hat{P}^3$

$$\delta_{p^+} O_\Delta = O_{\Delta+1}$$

TRANSLATION IS A SHIFT
IN CONFORMAL DIMENSION

CONFORMAL BASIS: SUMMARY

$\mathcal{O}_\Delta(z; X) \rightarrow$ PRIMARY STATE UNDER
LORENTZ TRANSFORMATIONS

SCALING DIMENSION: $\Delta = 1 + i\lambda \rightsquigarrow SL(2, \mathbb{C})$ PRINCIPAL
SERIES?

NOT SO CLEAR AS TRANSLATIONS RELATE \mathcal{O}_Δ WITH $\mathcal{O}_{\Delta+1}$

